

Thermochemical effects on hypersonic shock waves interacting with weak turbulence

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Huete: Spanish MCINN and BBVA Foundation (Leonardo Grant)

Urzay: US AFOSR and US DoE/NNSA



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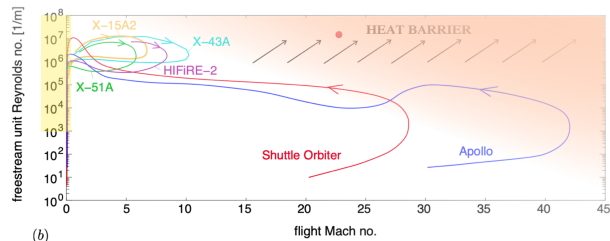
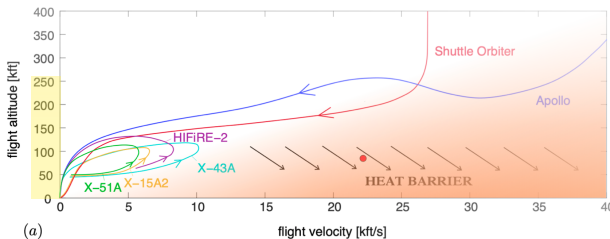
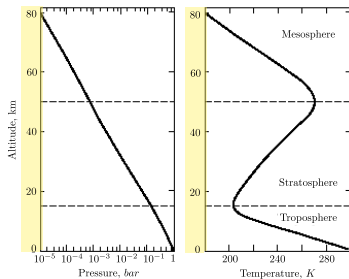
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Thermochemical effects on hypersonic shock waves interacting with weak turbulence

A. Cuadra (presenter)
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Urzay, J., & Di Renzo, M. (2021). Annual Research Briefs, Center for Turbulence Research, 7-32.

$$\frac{\Delta p}{p} \sim \frac{\Delta \rho}{\rho} = O(10^5), \quad \frac{\Delta T}{T} = O(1).$$

In hypersonic flight near the ground, the Reynolds number becomes large because of the comparatively larger densities

$$\frac{\Delta Re}{Re} = O(10^5).$$

1 Motivation

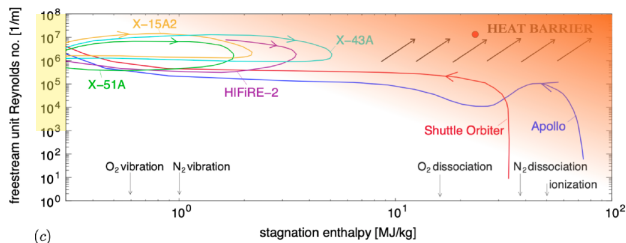
Effects of dissociation and vibrational excitation on the mean post-shock quantities

LIA of turbulence interacting with hypersonic shocks

Extension for real air mixture

Conclusion





Urzay, J., & Di Renzo, M. (2021). Annual Research Briefs, Center for Turbulence Research, 7-32.

Hypersonic flight **at low altitudes** is characterized by:

- High free-stream Mach numbers $Ma \geq 5$
- High free-stream and post-shock unit Reynolds numbers $Re \sim 10^7 - 10^9 \text{ m}^{-1}$
- High stagnation enthalpies $h_0 \sim 5 - 30 \text{ MJ/kg}$
- Small mean free paths $\lambda \sim 0.1 \text{ }\mu\text{m}$
- large normal Mach numbers
- turbulent boundary layers
- much higher than the vibrational specific energies of O_2 and N_2
- short vibrational relaxation distances

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2 Motivation

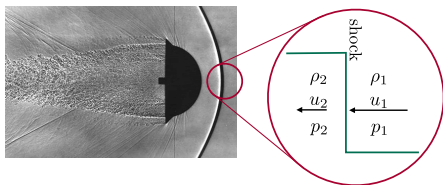
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Integral conservation equations across shock waves in **dissociating** gases

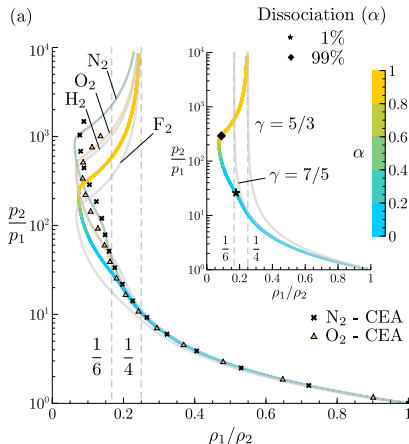
$$\rho_1 u_1 = \rho_2 u_2,$$

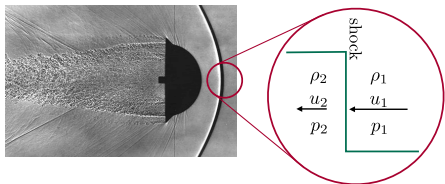
$$p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2,$$

$$e_1 + p_1/\rho_1 + u_1^2/2 = e_2 + p_2/\rho_2 + u_2^2/2 + q_d,$$

Upstream flow is sufficiently cold to be approximated as

$$e_1 = (5/2)R_{g,A_2}T_1, \quad p_1 = \rho_1 R_{g,A_2}T_1.$$





Integral conservation equations across shock waves in **dissociating** gases

$$\rho_1 u_1 = \rho_2 u_2,$$

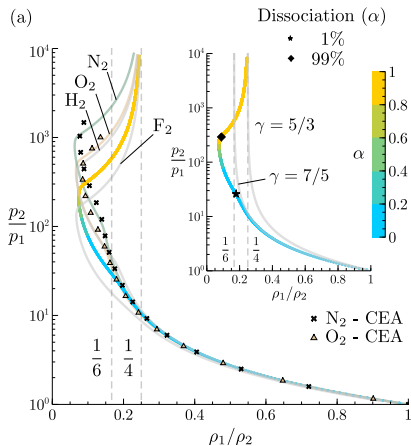
$$p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2,$$

$$e_1 + p_1/\rho_1 + u_1^2/2 = e_2 + p_2/\rho_2 + u_2^2/2 + q_d,$$

$$p_2 = \rho_2 R_{g,A_2} T_2 (1 + \alpha), \quad q_d = \alpha R_{g,A_2} \Theta_d,$$

$$e_2 = R_{g,A_2} T_2 \left[3\alpha + (1 - \alpha) \left(\frac{5}{2} + \frac{\Theta_v/T_2}{e^{\Theta_v/T_2} - 1} \right) \right],$$

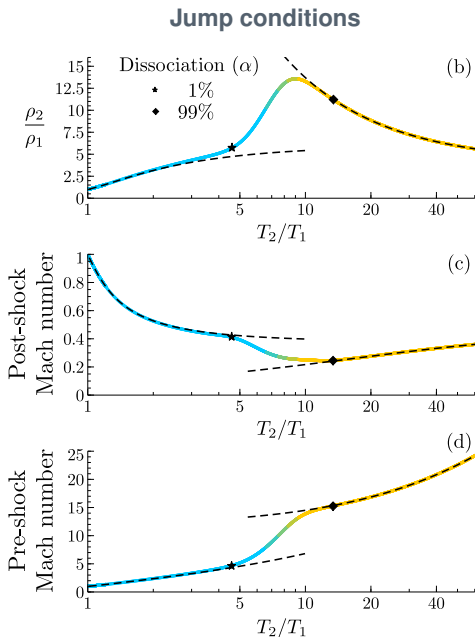
$$\frac{\alpha^2}{1 - \alpha} = Gm\Theta_r \left(\frac{\pi m k_B}{\hbar^2} \right)^{3/2} \frac{\sqrt{T_2}}{\rho_2} e^{-\frac{\Theta_d}{T_2}} \left(1 - e^{-\frac{\Theta_v}{T_2}} \right),$$



where α is the **degree of dissociation**, defined as the mass fraction of A atoms in the reaction $A_2 \rightleftharpoons A + A$, that must be solved with the aid of the chemical equilibrium condition.

Endothermicity due to dissociation and vibrational excitation does the following:

- increases the mean post-shock density
- decreases the mean post-shock velocity
- decreases the mean post-shock Mach number
- decreases the mean post-shock temperature

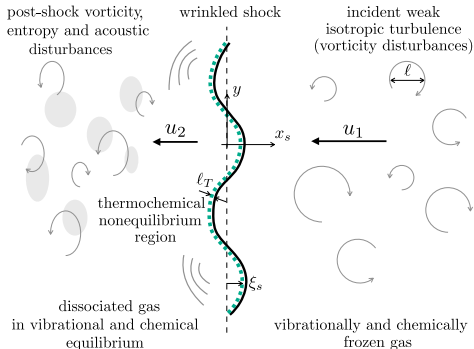


Considering:

- turbulence is comprised of small fluctuations,
- Kovaszny decomposition into vortical, entropic and acoustic modes,

we can solve this problem analytically by using

- linearized Rankine-Hugoniot relations,
- linearized Euler equations in the post-shock gas.



Limits of validity

Assumptions standard LIA:

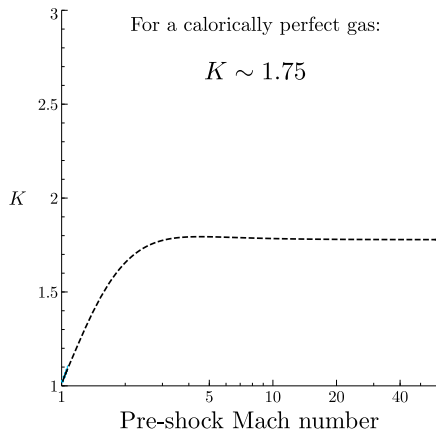
- $\text{rms}(u_\ell) \ll a_1$ and a_2 ,
- $\xi_s \ll l$,
- $l/u_\ell \ll l^2/\nu$.

With thermochemical effects:

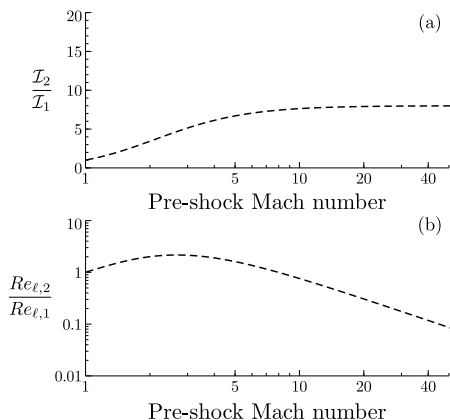
- $l_T \ll l$

For a given l , this condition becomes increasingly more accurate as the altitude decreases.

Turbulent Kinetic Energy

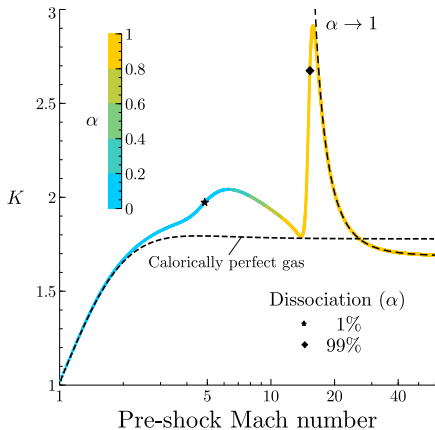


Turbulence intensity and Turbulent Reynolds number

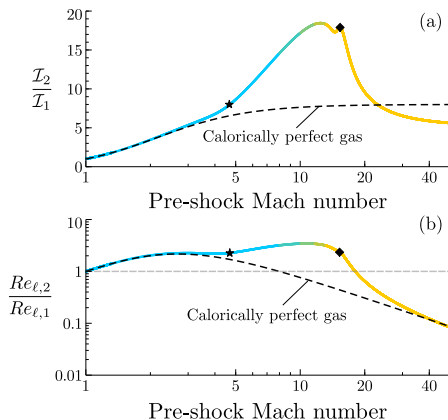


At hypervelocities ($Ma \gtrsim 10$), the calorically perfect gas approximation predicts a saturation in the amplification of kinetic energy and turbulence intensity, along with a decrease in the turbulent Reynolds number across the shock.

Turbulent Kinetic Energy

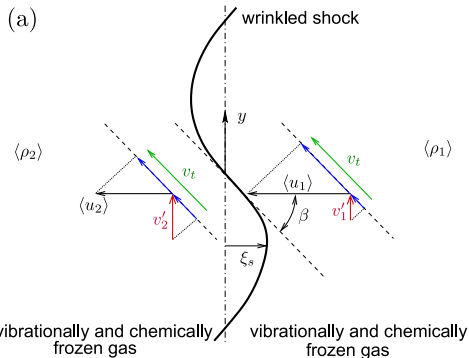


Turbulence intensity and Turbulent Reynolds number

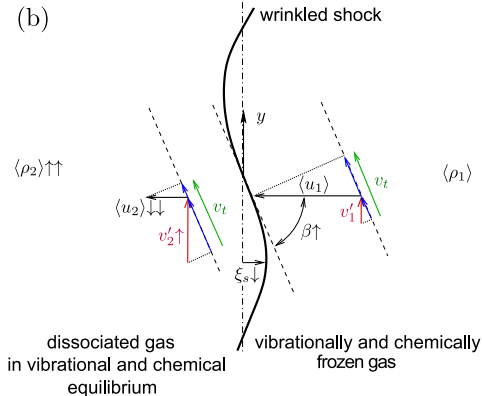


In contrast, the incorporation of dissociation and vibrational excitation predicts larger kinetic energy and turbulence intensity amplification rates, along with an increase in the turbulent Reynolds number across the shock.

Calorically perfect gas



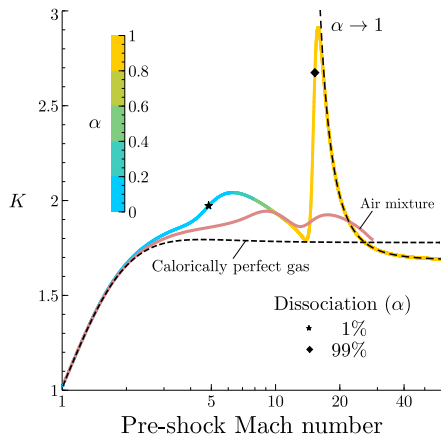
Vibrationally Excited, Dissociating Gas



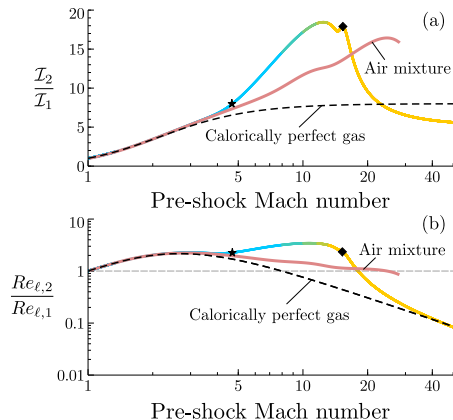
Conservation of tangential momentum dictates that the transverse velocity fluctuations should increase across the shock – these are larger at hypersonic velocities because of the associated larger post-shock densities induced by endothermic thermochemical effects.



Turbulent Kinetic Energy



Turbulence intensity and Turbulent Reynolds number



In case of a real air mixture (recombination in multi-species gas and electronic excitation) there is a significant decrease of the peaks values of TKE. However, the qualitative picture remains intact.

Key takeaways

- Significant departures from calorically perfect gas behaviour can be observed in the solution even at modest degrees of dissociation of 1%.
- **The amplification of TKE doubles** that observed in calorically perfect gases, with most of the content of TKE downstream in form of **vortical modes**.
- **The turbulent Reynolds number is amplified** across the shock at hypersonic Mach numbers in the presence of dissociation and vibrational excitation, as opposed to the attenuation observed in the calorically perfect case.
- Preliminary results show that real air mixture share qualitative results, but with lower amplitudes.
- **Thermochemical effects arising at hypersonic velocities appear to enhance turbulent fluctuations in the post-shock gas.**



Huete, C. et al. (2021). Thermochemical effects on hypersonic shock waves interacting with weak turbulence. *Physics of Fluids*, 33(8), 086111



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